

Microscopic rate coefficients:

$$k(E) = \frac{Q_{rot}^* W^*(E^*)}{Q_{rot} h \rho(E)} \quad (4.1.5)$$

$$W(E) \sim W_{cl}(E_{cl}) = \int_0^{E_{cl}} c_{vib-cl}^{(n_v)} \varepsilon_{cl}^{n_v-1} d\varepsilon_{cl} = c_{sum-cl}^{(n_v)} E_{cl}^{n_v} \quad (4.1.10)$$

⟨Whitten–Rabinovich approximation⟩

$$W(E) \sim c_{sum-cl}^{(n_v)} [E + (1 - \beta\omega)E_{zp}]^{n_v} \quad (4.1.11)$$

$$\beta = \frac{n_v - 1}{n_v} \frac{\langle \nu^2 \rangle}{\langle \nu \rangle^2}, \quad \begin{cases} \omega^{-1} = 5\varepsilon' + 2.73(\varepsilon')^{0.5} + 3.5100 & \text{for } \varepsilon' = 0.1 - 1.0 \\ \log_{10} \omega = -1.0506(\varepsilon')^{0.25} & \text{for } \varepsilon' = 1.0 - 8.0 \end{cases}, \quad \varepsilon' = E / E_{zp}$$

$$\rho(E) = \frac{dW}{dE} \sim c_{vib-cl}^{(n_v)} [E + (1 - \beta\omega)E_{zp}]^{n_v-1} \left(1 - \beta \frac{d\omega}{d\varepsilon'} \right) \quad (4.1.12)$$

Source list 1. A FORTRAN subroutine for the 'direct count'

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C-----
      SUBROUTINE DCOUNT(NUMV,JFRQ,NUMG,DENS,SUMS)
C      Direct count (Beyer-Swinehart algorithm)
C NUMV   : number of vibrators
C JFRQ() : vibrational frequencies (energy grain unit; integer)
C NUMG   : number of energy grains
C DENS() : density of states (subscript 1 corresponds to E=0)
C SUMS() : sum of states
C-----
      PARAMETER (MAXVIV=50,MAXGRN=80000)
      INTEGER NUMV,NUMG,JFRQ(MAXVIV),IV,IG,JFR,NSTT,JG
      DOUBLE PRECISION DENS(MAXGRN),SUMS(MAXGRN)
C      ----- Reset matrices
      DO IG=1,NUMG
        DENS(IG)=0.0D0
        SUMS(IG)=1.0D0
      ENDDO
      DENS(1)=1.0D0
C      ----- Beyer-Swinehart direct count
      DO IV=1,NUMV
        JFR=JFRQ(IV)
        NSTT=JFR+1
        DO IG=NSTT,NUMG
          JG=IG-JFR
          DENS(IG)=DENS(IG)+DENS(JG)
          SUMS(IG)=SUMS(IG)+SUMS(JG)
        ENDDO
      ENDDO
      RETURN
      END

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Problem-4.2

1) Evaluate $k(E)$ (eq. 4.1.5) for $\text{CH}_3\text{CH}_2\text{I} \rightarrow \text{CH}_2=\text{CH}_2 + \text{HI}$ at $E = 19000$ and 25000 cm^{-1} .

Compare the results of classical approximation [(2.3.5), (4.1.10)] and Whitten–Rabinovitch approximation [(4.1.11), (4.1.12)].

⟨ $\text{CH}_3\text{CH}_2\text{I}$ ⟩ vib. freq.: 3000(5), 1430(3), 1280(3), 1050(2), 940, 760, 500, 240, 200 cm^{-1}
 rot. const.: 0.0943, 0.101, 0.956 cm^{-1}

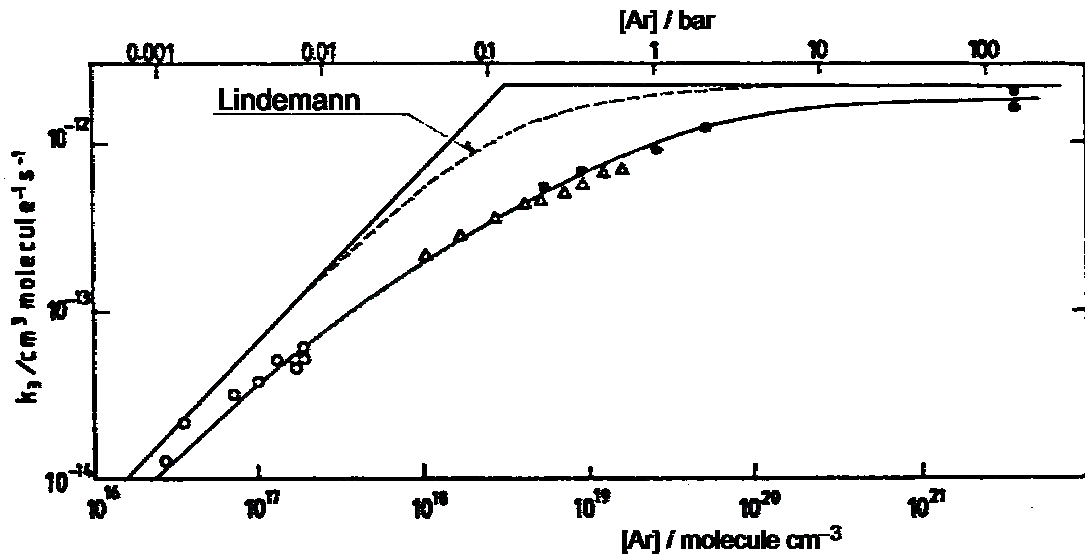
⟨TS⟩ vib. freq.: 3000(4), 1770, 1450, 1370(3), 1120, 1090, 1050, 930, 570, 330, 250, 170 cm^{-1}
 rot. const.: 0.0609, 0.0644, 0.774 cm^{-1}

Dissociation threshold energy $E_0 = 17800 \text{ cm}^{-1}$

2) [OPTION] Evaluate the density/sum of states by *direct count*, and compare with above results.

Unimolecular Reactions

$$\text{Lindemann: } k = \frac{k_0 k_\infty [M]}{k_0 [M] + k_\infty} \quad (4.2.4)$$

Fig. 4.4 Rate constant for $\text{CH}_3 + \text{O}_2$ recombination reaction

$$\text{Troe's formula: } k = \frac{k_0 k_\infty [M]}{k_0 [M] + k_\infty} F \quad (4.2.6)$$

RRKM theory

$$\text{RRKM: } k = \int_{E_0}^{\infty} k(E) g(E) dE = \frac{1}{Q_A} \int_{E_0}^{\infty} k(E) \frac{\rho(E) \exp(-E/k_B T)}{1 + k(E)/k_d [M]} dE \quad (4.3.3)$$

$$\text{H.P.L.: } k_\infty = \int_{E_0}^{\infty} k(E) F(E) dE = \frac{k_B T}{h} \frac{Q_{\text{rot}}^* Q^*}{Q_{\text{rot}} Q} \exp\left(-\frac{E_0}{k_B T}\right) = \text{TST} \quad (4.3.5)$$

$$\text{L.P.L.: } k_0 = k_d \int_{E_0}^{\infty} F(E) dE = \int_{E_0}^{\infty} k_a(E) dE \quad (4.3.6)$$

$$\text{Strong coll. RRKM: } k_d = k_d^{\text{SC}} = \Omega_{\text{A-M}}^{(2,2)*} \pi \sigma_{\text{A-M}}^2 \sqrt{\frac{8k_B T}{\pi \mu_{\text{A-B}}}} \quad (4.3.8)$$

$$\Omega_{\text{A-M}}^{(2,2)*} \sim [0.70 + 0.52 \log_{10}(k_B T / \varepsilon_{\text{A-M}})]^{-1}, \quad \sigma_{\text{A-M}} \sim (\sigma_A + \sigma_M)/2, \quad \varepsilon_{\text{A-M}} \sim \sqrt{\varepsilon_A \varepsilon_B}$$

$$\text{Weak coll. RRKM: } k_d = k_d^{\text{WC}} = \beta_c k_d^{\text{SC}} \quad (4.3.9)$$

Problem-4.4 [OPTION] _____

- 1) By using the weak collision corrected RRKM theory, calculate k_0 , k_∞ , and k at $[M] = [M]_c$ at 800 K for the reaction in the problem-4.2. Calculate ρ and W by the Whitten-Rabinovich approximation or by the direct count. Assume $\beta_c = 0.2$ and $M = \text{Ar}$, and use $\sigma(\text{CH}_3\text{CH}_2\text{I}) = 5.3 \text{ \AA}$, $\sigma(\text{Ar}) = 3.5 \text{ \AA}$, $\varepsilon(\text{CH}_3\text{CH}_2\text{I}) / k_B = 320 \text{ K}$, and $\varepsilon(\text{Ar}) / k_B = 93 \text{ K}$.
- 2) Plot the weak collision vibrational energy distribution, $g^{\text{WC}}(E)$, at $[M] = [M]_c$ against E , and compare it with the Boltzmann distribution, $F(E)$.